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ORBITALLY-DRIVEN MAGNETISM IN CORRELATED ELECTRON SYSTEMS

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Abstract. – Over the past decade, we have: (1) developed phenomenological theory for the behavior of “well-ordered” magnetic states of moderately delocalized light rare earth and actinide systems (characteristically obtaining unusual *anisotropic magnetism* in agreement with experiment); (2) developed theory and computational technique to synthesize first principles electronic structure information into that phenomenological theory to make it *materially predictive*. As discussed below, the resulting theory allows us to predict the triggering of an instability from unusual anisotropic, but well-ordered, magnetism to an unstable state. The unstable state can be either of a valence fluctuation type or of what probably is a heavy fermion type, and the detailed way in which these two types of instability is triggered differs.

1. Introduction

As the f-electron atomic shells are filled in the rare earths (4f's) and actinides (5f's), localization of the f electrons in space and energy occurs with increasing atomic number. For the rare earths, this transition to localized behavior occurs on increasing atomic number beyond that of neodymium. However, the degree of delocalization for the lighter rare earths (cerium, praseodymium, neodymium) depends on the detailed chemical environment (specific compound or alloy), and the degree and symptoms of partial delocalization are typically most obvious for cerium materials. For the actinides, the transition to localized f behavior occurs on going from plutonium to americium. Again, the degree of f delocalization for the light actinides (uranium, neptunium, plutonium) depends on the specific chemical environment. However, for the same chemical environment (same compound except for change of actinide), one expects considerably more delocalization for uranium than for plutonium. The important thing is that there is a *transitional regime with a range of partially delocalized behavior* between itinerant (band or bonding) and localized (correlated ionic) f-electron behavior. In recent years there has been much interest in understanding the solid-state behavior characteristic of this transitional regime. Magnetic behavior has played a central role in that interest, both as offering phenomena of interest in themselves, and as providing a diagnostic tool to understand how to treat electronic-structure-based behavior generally in the partially delocalized regime of behavior.

Over the past decade, we have: (1) developed phenomenological theory [1-10] for the behavior of “well-ordered” magnetic states of moderately delocalized light rare earth and actinide systems (characteristically obtaining unusual *anisotropic magnetism* in agreement with experiment); (2) developed theory and computational technique [11, 12] to synthesize first principles electronic structure information into that

phenomenological theory to make it *materially predictive*. As discussed below, the resulting theory allows us to predict the triggering of an instability from unusual anisotropic, but well-ordered, magnetism to an unstable state. The unstable state can be either of a valence fluctuation type or of what probably is a heavy fermion type, and the detailed way in which these two types of instability is triggered differs.

The key to understanding solid-state behavior in the transitional regime between localized and delocalized f-electron behavior is that the f-electrons belonging to one atom feel those belonging to another atom primarily through hybridization (mixing) with band electrons of non-f atomic parentage. There are two main themes involved in our theory for hybridization driven magnetic behavior and instabilities, corresponding to the theory of the magnetic phenomenology and to the synthesis of electronic structure information into that theory, respectively. The crucial point of the phenomenology is that partially delocalized f-electron systems have *orbitally-driven magnetism* [2]. The dominant magnetic coupling is between orbital moments *via* f electron-band electron hybridization (mixing). This gives a parallel channel to conventional spin-driven magnetism, and this orbitally-driven channel dominates the magnetic behavior as the f-electrons move from localized (gadolinium-like) behavior to partially delocalized (cerium-like) behavior. With regard to fitting band structure (material-specific) information into the phenomenological treatment *which includes onsite correlation*, it is important to develop the phenomenological (model Hamiltonian) theory and electronic structure theory (used to evaluate parameters of the phenomenological theory) in a unified way. Our phenomenological theory is a resonant scattering theory; [10, 13] and our electronic structure theory [11, 12] is a full-potential linearized muffin-tin orbital theory, which is based on a multiple scattering formalism, so that the two segments of theory fit together naturally.

Strongly anisotropic magnetic phenomena are characteristic of the hybridizing (moderately delocalized) regime of behavior [1-3]. The phenomenological (model Hamiltonian) theory as developed over the past decade has been very successful in explaining a large number of such phenomena in light rare earth and actinide systems. Among this variety of experimental behavior, there are three key experimental results that shaped or confirm the conceptual nature of the theory. The first of these is the seminal work of Vogt [14, 15] involving high field anisotropic magnetization experiments on CeSb and CeBi, giving extreme anisotropy that is the "wrong way around" (reversed) from that expected for single-ion crystal-field anisotropy; and that furthermore by diluting the cerium with an appropriate diluent (yttrium and/or lanthanum), one obtained the anisotropy behavior expected for single-ion crystal-field anisotropy. This showed the presence of an *extremely anisotropic two-ion interaction*. Second was the neutron inelastic scattering results of Birge-neau *et al.* [16] showing that the crystal-field splitting for CeSb and CeBi is an order of magnitude less than that expected by smooth extrapolation of the behavior of the heavier rare earth mononictides. Third, and much more recent [2, 3] is the recognition that the magnetic ordering temperatures of the light rare earth antimonides and bismuthides are much higher than would be expected in comparison to the behavior for the gadolinium and heavy rare earth compounds. The first of these results can be understood as coming from a *hybridization-mediated two-ion interaction*; the second can be understood on the basis of *hybridization-dressing of the crystal field*; and the third result can be understood on the basis of *hybridization giving rise to coupling between orbital ionic moments (i.e., orbitally driven magnetism)*.

2. Theory and methodology

Our understanding is based on the Anderson model Hamiltonian for a lattice of f-electron ions hybridizing with a band sea. The Hamiltonian then consists of terms for band and localized-f energies with a correlation energy cost (U) to add or subtract an f electron from the nominal f configuration of the ions, and a hybridization term which mixes band and f electrons. By performing a fourth-order Schrieffer-Wolff transformation [11], the hybridization is transformed to resonant scattering of band electrons off f electrons. This resonant scattering is equivalent to a virtual configuration transition for the f-electron ion. This gives rise both to a modification of the f-electron ionic behavior (hybridization dressing of the crystal field) and to two-ion interactions *via* correlated configuration transitions on two f-electron ions (two-ion interaction *via* cooperative hybridization with the band sea). The characteristic parameter giving the single-site band-f hybridiza-

tion (resonant scattering) strength is called $J(k, m, n, k'm'n')$ where k labels band states, m labels f states, and n labels the nominal f-ion configuration (e.g., f^1 for Ce^{3+}). J goes as the square of V_{km} , the band-f hybridization potential matrix element, and inversely as energy differences of the band (ε_k), f-electron (E_m) and correlation (nU) energies.

The two-ion hybridization-mediated exchange interaction is given essentially as a product of elements of J that are off-diagonal in k ; while the crystal-field dressing goes linearly as elements of J diagonal in k .

Two things have to be done to build a unified phenomenological/electronic structure theory. First, it is necessary to calculate the hybridization potential matrix elements on an absolute basis. Having both the phenomenological and the electronic structure theory in scattering language allows this to be done. Second, to the extent needed, it is necessary to simplify the band structure theory and the phenomenological theory sufficiently to see the physical nature of the dominant physical effect to be evaluated, and then "go backwards" to feed in the full complexity. Doing this is relatively easy for the crystal-field dressing, but is non-trivial for the two-ion interaction that is at the heart of the phenomenology, and involves reducing the size of the two-ion exchange matrix $E(m'_2, m_2; m'_1, m_1)$ to obtain initial physical understanding before inserting the full complexity of $E(m'_2, m_2; m'_1, m_1)$ into the absolute theory for the phenomenology.

To understand the second of these points, consider a Ce^{3+} lattice hybridizing with a band sea. In general for a given neighbor distance, the two-ion interaction matrix, $E(m_2, m_2; m_1, m_1)$ would have $6^4/4$ independent components (reduction by factor of 4 comes from time reversal symmetry). To gain physical feeling for what is going on depended on reducing the number of independent components per neighbor shell. (When we actually calculate, at present we can keep all $6^4/4$ components when we so desire.) If one has (approximate) axial symmetry about the two-ion axis, one conserves the component of angular momentum along the axis in transitions, and this reduces the number of independent parameters for a Ce^{3+} lattice to $6^2/4 = 9$. In mean field theory, hermiticity reduces this to 6, and if we take the large separation limit this is reduced further to a single component for each neighbor shell, namely that involving $m, m' = \pm 1/2$ with respect to the two-ion axis. These are the dominant components because they involve $m_l = 0$ components for the two ions, i.e., the cooperative hybridization effect tends to "suck out" ionic charge along the interionic axis. This then gives the picture of orbitally-driven anisotropic two-ion exchange interactions. Anisotropic, because for Ce^{3+} the ionic orbital moments tend to align perpendicular to the interionic axis because the hybridized disc-like ionic charge cloud wants to have the interionic

axis lying in the plane of the discs. (For light actinide ions involving more than one f electron this is modified by intraionic correlation.) The experimentally observed sensitivity to chemical environment can be easily understood when the close relationship between f-electron charge shaping and magnetic behavior is recognized.

In the past four years we have developed a computational scheme [11, 12] to perform a first principles evaluation of the parameters determining $J(k, m, n; k', m', n')$, i.e., the band, f, and correlation energies, and the hybridization potential matrix elements; and this has allowed us to put the phenomenological theory on an absolute basis *allowing us to predict changes in observable magnetic behavior on going from one material to another.*

This computational scheme involves a nonconventional electronic structure calculation. With hybridization suppressed, f-states are treated as resonant states in a solid-state environment and then localized at each iteration of a self-consistent electronic structure calculation. This provides the spatial dependence of self-consistently determined localized f states, band states and energies, and a self-consistently determined hybridization potential, thereby allowing determination of the hybridization potential matrix elements (mixing f and band electrons) which is the key ingredient of the theory. The f and correlation energies (in the sense of the model Hamiltonian) are determined from supercell calculations treating one cerium or actinide ion out of four as a defect differing from the nominal configuration by the addition or subtraction of an electron, and then using linear transition theory applied to the $4f^0$, $4f^1$ and $4f^2$ energies so calculated to obtain the f and correlation energies.

The parameters so calculated are then fed into the Hamiltonian obtained from the Schrieffer-Wolff transformation of the Anderson lattice Hamiltonian, which has the form of a resonant scattering interaction. The direct scattering gives the shift in crystal-field levels (hybridization dressing), and the exchange scattering gives a highly anisotropic two-ion anisotropic exchange interaction. Note that the theory provides two-ion interaction with every neighbor that is fully and absolutely determined both as to angular dependence and magnitude. We then use a mean field theory for this Hamiltonian to calculate the observable magnetic behavior.

It is important to realize that the mean field theory here involves quite a sophisticated procedure to look at the *partial* total energy associated with magnetic ordering, i.e., relative to the primary bonding total energy giving the cohesive energy of the solid. The interaction being represented by its time averaged behavior is a complex multipolar interaction. We have developed sophisticated numerical technique to find all reasonable possible metastable states, as well as the

ground state, at $T = 0$ and to follow these as temperature increases [1]. The iterative self-consistent procedure we have developed allows the system to seek whatever metastable states it wants. *The procedure does not require the existence of magnetic ordering and will tell us if state mixing (i.e., of ionic states quantized along directions that do not have to be collinear) leads to no magnetic ordering and a truly nonmagnetic state, i.e., without a paramagnetic moment.*

3. Results and discussion

The physical origin of the crystal-field dressing can be pictured rather easily. Unfilled band states of appropriate (Γ_7 and Γ_8) symmetry near the Fermi energy mix with crystal-field states of the same symmetry (Γ_7 or Γ_8 , respectively) and push them down. Our quantitative analysis [11] indicates that hybridization between f states and band states throughout the Brillouin zone, rather than just in the vicinity of high symmetry points, is important to understanding the magnitude of the relative shift in the crystal-field levels. The set of degenerate band states belonging to the star of a general point in the irreducible wedge of the Brillouin zone form a basis for a representation of the cubic group that, like the "regular" representation, contains the Γ_7 representation twice and the Γ_8 representation four times. (Here we are discussing cubic systems such as the NaCl-structure monpnictides and monochalcogenides.) Hence, with twice the number of states hybridizing with the cerium $4f$ Γ_8 quartet, we expect the effect on the Γ_8 quartet to be larger than the effect on the Γ_7 doublet by approximately a factor of two. Our results confirm this expectation. Both Γ_7 and Γ_8 crystal-field f states are lowered, with the effect on Γ_8 greater than the effect on Γ_7 by a factor of 2.45 for CeBi, 2.56 for CeSb, 2.16 for CeAs, and 2.09 for CeP.

Now we will discuss the use of our procedure to analyse the trend between materials in hybridization-induced effects, using the behavior of two quantities to provide a gross characterization of what is going on: (1) the calculated f-state resonance width, Γ , characterizing the strength of hybridization per appropriate band electron, and (2) the density of states at the Fermi energy, $N(\epsilon_F)$, characterizing the number and character (cerium d, anion p) of band states available for hybridization. We have now applied our ab initio treatment to understand [12] the strong change in nature of magnetic ordering on going from CeSb to CeTe (going from pnictide to chalcogenide means adding one anion valence p electron) both of which have NaCl-structure with a small decrease in lattice parameter. The key result is the way the ab initio absolute calculation signals the change from a stable well-ordered magnetic system to a system on the verge of an insta-

bility (probably to a heavy fermion state). In another application, to CeN, we can see how the existence of an unstable valence system (valence fluctuation system) is signaled [17]. Finally we will tabulate and discuss trends in hybridization strength more broadly.

3.1 CHANGE IN BEHAVIOR FROM CeSb TO CeTe. – CeSb is [14] an antiferromagnet with almost saturated (free-ion) ordered moments aligned along (100). The Néel temperature is 16.8 K and the ordered moment at low temperature is $2.06 \mu_B$. On the other hand, CeTe is [18] an antiferromagnet with very small ordered moment aligned along (111). The Néel temperature is only 2.4 K and the ordered moment at low temperature is only $0.20 \mu_B$.

On calculating the electronic structure for CeSb and CeTe treating the f state as a resonant state, we see that the addition of a p electron causes the Fermi level to rise into a region of high Ce 5d density of states. Thus the big change in behavior on going from CeSb to CeTe is that the bands dominating hybridization-induced properties in CeTe are largely due to f hybridization with band states largely derived from Ce 5d states [12]. The resonance width, Γ , only increases slightly on going from CeSb to CeTe, from 3.06 mRy to 3.21 mRy ; but $N(\epsilon_F)$, the density of states at the Fermi level increases from 4.38 Ry^{-1} to 15.56 Ry^{-1} where the increase is largely due to the d density of states increase from 1.46 Ry^{-1} to 11.43 Ry^{-1} . This implies a large net increase in hybridization because of the big increase in band states available for hybridization, and when we performed the sequence of calculations summarized in the previous section we find a strong change in the range, size and signs of the two-ion exchange parameters with the various neighbor ions on going from CeSb to CeTe. If we treat the two-ion exchange matrix at the highest level of approximation, including the large ionic separation limit, as described in the previous section there is only one two-ion exchange parameter per neighbor shell. A calculation at this level of approximation gives a free-energy per Ce^{3+} ion associated with magnetic ordering of -2.6 K for CeSb, but only -0.3 K for CeTe. This means that CeSb is a well-magnetically-ordered system, while CeTe is only weakly-ordered. Indeed the calculations give an ordered moment close to saturation for CeSb and predict a ratio of Néel temperatures of about 10, and also predict the correct change in direction of ordered moment. The degree of approximation of the two-ion exchange matrix can be improved by eliminating the large ionic separation limit approximation, i.e., level of treatment where these are now 6 calculated exchange parameters per neighbor shell (remember everything used in the calculations is calculated wholly ab initio). On doing this, the free energy associated with magnetic ordering for CeSb increases further in magnitude to -3.8 K , i.e., the well-

ordered state is "confirmed" by improving the degree of approximation. However, for CeTe, there is now no stable magnetically-ordered state at $T = 0$. *Thus the increase in band-f hybridization places CeTe at the borderline between magnetic and nonmagnetic behavior in contrast to almost fully-ordered CeSb.*

3.2 PREDICTION OF VALENCE FLUCTUATION BEHAVIOR FOR CeN. – As discussed in the previous section, the f-state and correlation energies, E_f and U , are obtained by performing a sequence of three supercell calculations (unit cell equal to four normal unit cells) treating one cerium as a defect having f^0 or f^2 configuration (as well as a "control" supercell calculation where the defect site has the same f^1 configuration as the other sites). The values of E_f and U are then obtained from the f^0 , f^1 and f^2 energy levels by the use of linear transition theory. For the heavier cerium mononpnictides this procedure gives the $4f^1$ level about 0.5 eV below the Fermi level; and E_f is about 3 eV below the Fermi level [17], with U equal to about 6 eV . The resonance width, Γ , for the heavier cerium mononpnictides varies between a low value of about 3 mRy for CeSb to a high value of about 5 mRy for CeP. For CeN, the value of Γ increases [17] to 21.2 mRy ; as discussed below this is associated with a large increase of f hybridization with Ce 5d band states. The calculated $4f^1$ level moves to 2.1 eV above the Fermi level while the $4f^0$ level is 4.8 eV below the Fermi level. When we recognize that the Fermi level is an infinite sink of electrons, this movement of $4f^1$ above the Fermi level indicates a valence instability, a fluctuation between f^1 and f^0 ; and CeN is experimentally observed [19] to be an unstable valence system.

3.3 TRENDS IN HYBRIDIZATION INDICATED BY RESONANCE WIDTH. – As shown in table I, we can use the value of calculated resonance width, Γ , as one way of characterizing the trend between materials in strength of hybridization. (Remember that band densities of states also indicate changes in hybridization. Also we

Table I. – Calculated resonance width, Γ , for cerium and plutonium mononpnictides and monochalcogenides.

	$\Gamma \text{ (mRy)}$		$\Gamma \text{ (mRy)}$
CeN	21.22		
CeP	4.99		
CeAs	4.28		
CeSb	3.06	CeTe	3.21
CeBi	4.23		
PuAs	19.34		
PuSb	10.49	PuTe	17.86
PuBi	10.55		

expect a decreased calculated correlation energy U as another result of increased hybridization.)

Note the overall increase in Γ on going between corresponding cerium and plutonium compounds. This corresponds to greater f delocalization, hence increased hybridization. We expect another significant increase on going to uranium.

For the cerium monpnictides note the minimum in Γ at CeSb. For the lighter cerium monpnictides (CeN-CeSb), the lattice parameter is of a size that just accommodates cerium and pnictogen atoms of "typical solid-state size"; while for CeBi the lattice parameter is too small to do this. Thus the increase of hybridization on going up the pnictogen column from the antimonide has a simple explanation in increased f hybridization with 5d electrons on other cerium sites as the cerium-cerium separation decreases with decreasing anion size; the decrease on going to the nitride is large, giving a large increase in hybridization. On the other hand, the increase in Γ on going from CeSb to CeBi corresponds to the increased Ce f hybridization with anion valence p electrons as the cerium and bismuth draw closer together than would be expected from simple solid-state atomic size considerations.

It is well to remember that table I does not show the large increases in hybridization that can be associated with density of states effects. As discussed in subsection 3.1 above, such effects drive CeTe to the edge of magnetically unstable behavior. *It is very interesting that PuTe is the most anomalous material in table I, probably being a heavy fermion system* [20-22]. In PuTe we have a large increase in resonance with from PuSb, and in addition we expect a large increase in density of states of the same sort as occurs between CeSb and CeTe.

In conclusion, we believe we have developed a computational technique synthesizing model hamiltonian and electronic structure theory that reliably allows us to understand and predict which light rare earth and light actinide materials should have hybridization-driven magnetic and valence instabilities.

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- [1] Cooper, B. R., Siemann, R., Yang, D., Thayamballi, P. and Banerjea, A., *Handbook on the Physics and Chemistry of the Actinides*, Eds. A. J. Freeman and G. H. Lander (North-Holland, Amsterdam Vol. 2 (1985) pp. 435-500.
- [2] Cooper, B. R., *J. Less-Common Met.* **133** (1987) 31.
- [3] Cooper, B. R., Hu, G. J., Kioussis, N. and Wills, J. M., *J. Magn. Magn. Mater.* **63** (1987) 121.
- [4] Cooper, B. R. and Siemann, R. S., *J. Appl. Phys.* **50** (1979) 1991; Siemann, R. and Cooper, B. R., *Phys. Rev. Lett.* **44** (1980) 1015.
- [5] Thayamballi, P. and Cooper, B. R., *J. Appl. Phys.* **55** (1984) 1829.
- [6] Banerjea, A. and Cooper, B. R., *Phys. Rev. B* **34** (1986) 1607.
- [7] Kioussis, N. and Cooper, B. R., *Phys. Rev. B* **34** (1986) 3261.
- [8] Kioussis, N., Cooper, B. R. and Banerjea, A., *J. Appl. Phys.* **61** (1987) 3388; *Phys. Rev. B* **38** (1988) 9132.
- [9] Hu, G. J. and Cooper, B. R., *J. Appl. Phys.* **63** (1988) 3826; Thayamballi, P. and Cooper, B. R., *Phys. Rev. B* **31** (1985) 5911.
- [10] Hu, G. J., Kioussis, N., Banerjea, A. and Cooper, B. R., *Phys. Rev. B* **38** (1988) 2639; Hu, G. J. and Cooper, B. R., *Phys. Rev. B* **38** (1988) 9127.
- [11] Wills, J. M. and Cooper, B. R., *Phys. Rev. B* **36** (1987) 3809.
- [12] Kioussis, N., Cooper, B. R. and Wills, J. M., *J. Appl. Phys.* **63** (1988) 3683.
- [13] Coqblin, B. and Schrieffer, J. R., *Phys. Rev.* **185** (1969) 847.
- [14] Busch, G. and Vogt, O., *Phys. Lett. A* (1967) 449.
- [15] Cooper, B. R. and Vogt, O., *J. Phys. Colloq. France* **31** (1970) C1-1026.
- [16] Birgeneau, R. J., Bucher, E., Maita, J. P., Passell, L. and Turberfield, K. C., *Phys. Rev. B* **8** (1973) 5345.
- [17] Sheng, Q.-G., Cooper, B. R., Kioussis, N. and Wills, J. M., *Bull. Am. Phys. Soc.* **33** (1988) 364.
- [18] Hulliger, F., Natterer, B. and Ott, H. R., *J. magn. Magn. Mater.* **8** (1978) 87; Ott, H. R., Kjems, J. K. and Hulliger, F., *Phys. Rev. Lett.* **42** (1979) 1378.
- [19] Baer Y. and Zurcher, Ch., *Phys. Rev. Lett.* **39** (1977) 956.
- [20] Mattenberger, K., Vogt, O., Spirlet, J. C. and Rebizant, J., *J. Less-Common Met.* **121** (1986) 285.
- [21] Collard, J. M., Blaise, A., Fournier, J. M. and Charvillat, J. P., *J. Less-Common Met.* **121** (1986) 223.
- [22] Therond, P. G., Blaise, A., Fournier, J. M., Rossat-Mignod, J., Spirlet, J. C., Rebizant, J. and Vogt, O., *J. Magn. Magn. Mater.* **63** (1987) 142.